

Phase Transitions in Quantum Spin Systems with Long-Range Interactions: A Nonlinear Variational Inequality Approach

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Article History:

Received: 10-01-2025

Revised: 27-02-2025

Accepted: 07-03-2025

Abstract

This paper presents a novel mathematical framework for modeling phase transitions in quantum spin systems with long-range interactions using nonlinear variational inequalities (NVIs). Traditional methods such as mean-field theory and Landau-Ginzburg models often fall short in capturing spatial inhomogeneities, critical phenomena, and constraints inherent in realistic spin systems. Our approach addresses these limitations by reformulating the energy-based governing equations as an NVI problem defined over a suitable convex set of physically admissible spin configurations. We incorporate both Ising and Heisenberg-type quantum lattice models with power-law decaying interactions to accurately represent nonlocal effects. This research provides rigorous mathematical preliminaries, including operator-theoretic foundations in Hilbert and Banach spaces, and prove existence and uniqueness results under monotonicity and coercivity assumptions. Analytical insights include bifurcation analysis, critical point theory, and regularity of solutions, offering a robust theoretical foundation for understanding phase transitions. A finite element–Galerkin discretization is developed to simulate the variational system numerically. Simulation results demonstrate the emergence of ordered phases, spin textures, and domain formation, consistent with known physical behavior. Comparisons with classical models reveal the advantages of the NVI framework in capturing phase boundaries and metastable states with higher fidelity. The proposed framework also opens pathways for extending the analysis to quantum entanglement, topological order, and non-equilibrium phenomena. This research thus contributes both to the mathematical theory of variational inequalities and to the physical modeling of quantum spin systems, providing tools for deeper exploration of complex condensed matter phenomena.

Keywords: Quantum spin systems, phase transitions, nonlinear variational inequalities, long-range interactions, critical phenomena

1. INTRODUCTION

Phase transitions in quantum spin systems have been a subject of enduring interest in both theoretical physics and mathematical modeling. These systems exhibit rich behavior, including quantum criticality, topological order, and entanglement-driven phenomena, especially in low-dimensional materials and strongly correlated regimes. As quantum materials and spintronic devices continue to evolve, understanding the nature and mathematical structure of phase transitions has become more pressing (Sachdev, 2019; Frérot & Roscilde, 2018). Long-range interactions such as dipolar, Coulombic, or Rydberg-mediated couplings—are increasingly central to quantum spin models, offering mechanisms for exotic phase behavior and nonlocal correlations (Defenu et al., 2021). These interactions, decaying typically with a power-law $r^{-\alpha}$, significantly alter the critical exponents and universality classes of the system, posing new challenges for both physical interpretation and mathematical analysis (Maghrebi, 2017; Gori et al., 2022). Unlike nearest-neighbor systems, long-range interactions require

new analytical tools that can handle nonlocal operators and singular kernels, particularly when dealing with spin networks in infinite-volume or continuum limits.

Traditional approaches to phase transitions such as mean-field theory, renormalization group (RG) techniques, or Monte Carlo methods often face limitations in rigor and resolution when addressing strongly nonlinear, infinite-dimensional, or nonconvex problem domains (Henkel et al., 2023). Moreover, such methods may obscure the mathematical structure underlying solution spaces and stability near critical points. In this context, nonlinear variational inequalities (NVI) provide a powerful alternative framework, offering a rigorous formulation of equilibrium states subject to nonlinear constraints and dynamic energy landscapes. NVIs are well-suited for handling constraint-bound evolution (fixed spin magnitudes or bounded fields), and enable analytical characterization of solution existence, uniqueness, and bifurcation behavior central to understanding phase transitions (Faccioli et al., 2020). Despite their widespread use in mechanics and optimization, their application to quantum spin systems with long-range interactions remains underexplored.

This paper aims to develop a novel NVI-based formulation for modeling phase transitions in quantum spin systems with long-range interactions. By embedding the system's energy functional into a constrained variational inequality setting, we aim to derive conditions for criticality, explore solution regularity, and analyze phase structure transitions through a mathematically rigorous lens. The primary objective of this study is to develop a rigorous mathematical framework based on nonlinear variational inequalities (NVI) to analyze phase transitions in quantum spin systems with long-range interactions. While phase transitions in such systems have traditionally been studied using numerical simulations, renormalization group analysis, and mean-field approximations, these methods often fall short in capturing the rich structural and dynamical behavior introduced by long-range couplings—especially in infinite-dimensional settings or under physical constraints. This work seeks to bridge that gap by offering an analytical and variational formulation that accommodates the complexities of long-range interactions.

2. MATHEMATICAL PRELIMINARIES

To formalize the variational inequality model for phase transitions in quantum spin systems, we begin by establishing the functional analytic setting and the foundational tools required for the subsequent analysis.

We consider a separable Hilbert space H , typically $L^2(\Omega)$, where $\Omega \subseteq \mathbb{R}^d$ represents the spatial domain of the spin system. The norm and inner product on H are denoted by $\|\cdot\|$ and $\langle \cdot, \cdot \rangle$, respectively. The choice of a Hilbert space is particularly appropriate due to its reflexivity and the well-defined dual pairing structure, which are essential in the variational inequality framework.

Let $K \subseteq H$ be a nonempty, closed, and convex set representing the physically admissible configurations of the spin field—often constrained by spin magnitude or boundary conditions. The variational inequality problem is to find $u \in K$ such that:

$$\langle A(u), v - u \rangle \geq \langle f, v - u \rangle \quad \forall v \in K$$

where $A: H \rightarrow H^*$ is a nonlinear operator derived from the system's energy functional, and $f \in H^*$ represents external forcing or field terms.

We assume that A satisfies monotonicity and coercivity, two critical properties that ensure solvability of the inequality. An operator A is monotone if

$$\langle A(u) - A(v), u - v \rangle \geq 0 \quad \forall u, v \in H, \text{ and coercive if}$$

$$\langle A(u), u \rangle / \|u\| \rightarrow \infty \text{ as } \|u\| \rightarrow \infty.$$

Two classical results underlie the theoretical guarantees for solution existence:

- The Browder–Minty Theorem ensures that a coercive, hemicontinuous, and monotone operator A on a reflexive Banach space admits a solution to the corresponding variational inequality.
- The Lions–Stampacchia Theorem guarantees the existence of a unique solution to variational inequalities in Hilbert spaces when A is linear and coercive, which can be extended under nonlinear conditions.

In the context of quantum spin systems, energy functionals often take the form:

$$E[u] = 1/2 \iint J(x,y)(u(x) - u(y))^2 dx dy + \int V(u(x)) dx$$

where $J(x, y) \sim |x - y|^{-\alpha}$ encodes long-range interactions and $V(u)$ represents local potential terms. Such nonlocal terms naturally give rise to non-symmetric, nonlinear operators, making the variational inequality approach both necessary and robust. This mathematical framework sets the stage for reformulating the phase transition problem into a nonlinear variational inequality, enabling rigorous investigation of equilibrium states and bifurcation phenomena in complex quantum systems.

3. PHYSICAL MODEL FORMULATION

Quantum spin systems are typically modeled on a discrete lattice $\Lambda \subset \mathbb{Z}^d$, where each site $i \in \Lambda$ hosts a quantum spin variable σ_i . Two prototypical models for describing such systems are the Ising model and the Heisenberg model, distinguished by the degrees of freedom and symmetry of spin interactions. In the quantum Ising model, the spin variables take values $\sigma_i \in \{-1, +1\}$, and interactions are scalar-valued. The Hamiltonian for a long-range Ising system in transverse field is given by:

$$H = -\sum_{\{i \neq j\}} J_{\{ij\}} \sigma_i^z \sigma_j^z - h \sum_i \sigma_i^x,$$

where $J_{\{ij\}} \sim |i - j|^{-\alpha}$ is a long-range coupling term with decay exponent $\alpha > 0$, σ_i^z and σ_i^x are Pauli matrices, and h is the transverse field strength. For a more general vector-valued setting, the Heisenberg model considers spin variables $\sigma_i \in \mathbb{R}^3$, constrained to lie on the unit sphere $|\sigma_i| = 1$. The Hamiltonian becomes:

$$H = -\sum_{\{i \neq j\}} J_{\{ij\}} \sigma_i \cdot \sigma_j,$$

capturing the isotropic interaction between neighboring spins over long distances. **Continuum Limit and Energy Functional.** To transition from the discrete lattice to a continuum model, we consider a macroscopic spin density field $u(x): \Omega \rightarrow \mathbb{R}$ (or \mathbb{R}^3 for vector spins), where $\Omega \subset \mathbb{R}^d$ is a bounded domain representing the spatial extent of the material. The corresponding nonlocal energy functional is given by:

$$E[u] = (1/2) \iint J(x, y)(u(x) - u(y))^2 dx dy + \int V(u(x)) dx,$$

where:

- $J(x, y) = 1 / |x - y|^\alpha$, $\alpha \in (0, d)$, governs the strength of long-range interactions.
- $V(u)$ is a local potential term

This formulation captures nonlocal interactions characteristic of physical systems with algebraically decaying couplings, as observed in dipolar and Rydberg-interacting spin arrays.

Governing Equations

The equilibrium configurations of the spin system correspond to the minimizers of the energy functional $E[u]$, subject to admissibility constraints. The associated Euler–Lagrange equation, derived formally, reads:

$$\int J(x, y)(u(x) - u(y)) dy + V'(u(x)) = 0, \quad x \in \Omega.$$

This nonlocal integral equation can be interpreted as a balance between interaction forces and local field effects. In practice, spin configurations may also be constrained by physical bounds (e.g., $|u(x)| \leq 1$), in which case the equilibrium problem is more appropriately formulated as a nonlinear variational inequality. Under this formulation, the goal is to find $u \in K \subset H$, such that:

$$\langle A(u), v - u \rangle \geq 0 \quad \forall v \in K,$$

where $A(u)$ is the nonlinear operator associated with the first variation of $E[u]$, incorporating both nonlocal and potential terms. This variational framework allows for a mathematically rigorous analysis of phase behavior in long-range spin systems, while maintaining fidelity to the underlying quantum physics. The next section develops the full variational inequality model and examines its analytic properties.

4. VARIATIONAL INEQUALITY FRAMEWORK

To rigorously analyze the equilibrium behavior of quantum spin systems with long-range interactions, we reformulate the energy minimization problem as a nonlinear variational inequality (NVI) problem. This approach not only encapsulates the nonlocal nature of the interactions but also integrates essential physical constraints on spin configurations.

4.1 Reformulation into a Variational Inequality

Let $H = L^2(\Omega)$ denote the Hilbert space of square-integrable spin fields over a bounded domain $\Omega \subset \mathbb{R}^d$. Define the energy functional:

$$E[u] = (1/2) \int_{\Omega} \int_{\Omega} J(x, y)(u(x) - u(y))^2 dx dy + \int_{\Omega} V(u(x)) dx,$$

where:

- $J(x, y) = 1 / |x - y|^\alpha$, $\alpha \in (0, d)$, characterizes the long-range interactions,
- $V(u)$ is a nonlinear potential, such as a double-well potential modeling bistable phases.

The formal Euler–Lagrange equation associated with the stationary points of $E[u]$ yields the nonlocal equation:

$$\int_{\Omega} J(x, y)(u(x) - u(y)) dy + V'(u(x)) = 0, \text{ for } x \in \Omega.$$

With constraints like $u(x) \in [-1, 1]$, the problem becomes more naturally expressed as a nonlinear variational inequality: find $u \in K \subset H$ such that

$$\langle A(u), v - u \rangle \geq 0 \quad \forall v \in K,$$

where $A: H \rightarrow H^*$ is a nonlinear operator defined by:

$$\langle A(u), v \rangle = \int_{\Omega} \int_{\Omega} J(x, y)(u(x) - u(y))(v(x) - v(y)) dx dy + \int_{\Omega} V'(u(x)) v(x) dx.$$

4.2 Feasible Set

The feasible set $K \subset H$ encodes the physical admissibility of spin configurations:

$$K = \{ u \in H \mid |u(x)| \leq 1 \text{ almost everywhere in } \Omega \},$$

which is nonempty, convex, and closed in H .

4.3 Operator Properties

Key properties of the operator A include:

- Monotonicity:

$$\langle A(u) - A(v), u - v \rangle \geq 0 \text{ for all } u, v \in K.$$

- Coercivity:

$$\langle A(u), u \rangle \geq c \|u\|^2 - C \text{ for constants } c > 0, C \geq 0.$$

- Hemicontinuity: The map $t \mapsto \langle A(u + tv), w \rangle$ is continuous for fixed $u, v, w \in H$.

4.4 Existence and Uniqueness of Solutions

Assuming K is convex, closed, and nonempty in H , and A is coercive, monotone, and hemicontinuous, the Browder–Minty theorem guarantees existence of at least one solution $u^* \in K$ such that

$$\langle A(u^*), v - u^* \rangle \geq 0 \quad \forall v \in K.$$

If A is strictly monotone (e.g., $V''(u) \geq \delta > 0$ for $\delta > 0$), uniqueness of the solution is ensured. This NVI formulation thus provides both theoretical robustness and physical fidelity for modeling phase transitions in quantum spin systems.

5. Analytical Results

This section presents a detailed analytical exploration of the nonlinear variational inequality (NVI) model formulated for quantum spin systems with long-range interactions. The focus lies on identifying phase transition phenomena through bifurcation analysis, understanding how system parameters affect solution behavior, and comparing this framework to classical mean-field models.

5.1 Bifurcation and Critical Points

In the context of quantum spin systems, phase transitions are often characterized by bifurcations in the set of equilibrium solutions. Let the control parameter λ represent quantities such as interaction strength or inverse temperature. The NVI formulation becomes:

$$\langle A_\lambda(u), v-u \rangle \geq 0 \quad \forall v \in K$$

where A_λ is a nonlinear, parameter-dependent operator. Bifurcation theory indicates that for certain critical values λ_c the system undergoes qualitative changes in solution structure, such as:

- **Symmetry-breaking bifurcation**, where a uniform spin state loses stability and gives rise to distinct ordered phases.
- **Subcritical and supercritical bifurcations**, depending on the nature of the energy landscape and potential $V(u)$.

Linearized stability analysis near bifurcation points involves studying the Fréchet derivative $DuA_\lambda(u)$. The loss of invertibility of this linear operator signals the onset of bifurcation.

5.2 Dependence on System Parameters

The behavior of solutions is significantly influenced by system parameters, particularly:

- **Interaction decay exponent α** : Smaller values (i.e., slower decay of interactions) lead to enhanced coherence across the lattice and lower critical temperatures for phase transitions.
- **Inverse temperature β** : Higher β (i.e., lower temperatures) amplify the nonlinearity of the potential $V(u)$, creating sharper phase boundaries.
- **External field or anisotropy terms**, if present, introduce asymmetry and can tilt the bifurcation behavior in favor of one phase.

Parametric studies help map regions in the (α, β) -space where transitions between disordered and ordered spin phases are observed.

5.3 Solution Regularity and Stability

Assuming $J(x, y)$ and $V(u)$ are sufficiently smooth and satisfy structural conditions, the solution $u \in K$ belongs to fractional Sobolev spaces $H^s(\Omega)$ with $s \in (0, 1)$ due to the nonlocal nature of the operator. Specific results include:

- **Hölder continuity** of solutions under appropriate assumptions on kernel regularity and boundary conditions.
- **Energy stability**, assessed via the second variation of the energy functional:

$$\delta^2 E[u](v, v) = \iint J(x, y) (v(x) - v(y))^2 dx dy + \int V''(u(x)) v(x)^2 dx$$

A positive definite second variation implies that the solution u is a local minimizer and hence stable under small perturbations.

5.4 Comparison with Classical Mean-Field Models

Classical mean-field models reduce interactions to spatial averages, leading to simplified local equations like:

$$V'(u) = \mu u$$

where μ is an effective interaction coefficient. However, such models:

- Neglect spatial correlations and fluctuations,
- Fail to capture realistic critical exponents in low-dimensional systems,

- Do not accommodate inhomogeneities or metastable states.

6. Numerical Methods and Simulations

This section outlines the numerical framework developed to approximate and simulate the nonlinear variational inequality (NVI) arising in quantum spin systems with long-range interactions. The aim is to construct stable and convergent schemes capable of capturing complex spin behavior, including phase transitions, and to validate the model against known benchmarks.

6.1 Discretization Scheme

Given the variational inequality formulation in a function space setting, we employ a Galerkin-type finite element method (FEM) for discretization. Let $\Omega \subset \mathbb{R}^d$ be the spatial domain (typically a bounded lattice subset), and define a conforming finite-dimensional subspace $V_h \subset H^s(\Omega)$ (e.g., piecewise linear elements over a triangulation).

The discrete problem becomes: find $u_h \in K_h \subset V_h$ such that

$$\langle A_h(u_h), v_h - u_h \rangle \geq 0 \quad \forall v_h \in K_h$$

where A_h is the discretized nonlinear operator incorporating long-range interactions via quadrature approximations of the nonlocal integral kernel $J(x, y)$. The feasible set K_h enforces box constraints or admissibility conditions (e.g., spin values within $[-1, 1]$).

6.2 Simulation Results

To explore the physical behavior of the system under the NVI model, we simulate the model on a 2D periodic lattice using long-range interaction kernels of the form $J(x, y) \sim |x - y|^{-\alpha}$. The simulations produce the following outcomes:

- **Phase Diagrams:** Computed over a range of parameters (α, β) , revealing ordered and disordered regions. The critical curve delineates phase transition boundaries.
- **Spin Configuration Snapshots:** Visualizations of spin field $u(x)$ for various α, β , showing symmetry-breaking, domain formation, and coexistence of phases.
- **Energy Evolution:** Monotonic decay of the system energy functional during iterations confirms numerical stability.

6.4 Validation with Known or Simplified Cases

To verify the accuracy and robustness of the numerical method, simulations are compared with:

- **Mean-Field Solutions:** In the limit of infinite-range interaction ($\alpha \rightarrow 0$), the system approximates a classical mean-field model. The numerical solutions agree with analytical results in this limit.
- **1D Analytical Cases:** For specific 1D kernels $J(x, y)$ and potential $V(u)$ analytical solutions are available. Numerical results show excellent agreement with these benchmark cases.
- **Conservation Tests:** Numerical experiments conserve global magnetization (where applicable) and respect symmetry properties of the spin interaction kernel.

7. DISCUSSION

The adoption of a nonlinear variational inequality (NVI) framework for modeling phase transitions in quantum spin systems with long-range interactions offers both theoretical and practical advancements in understanding the critical behavior of complex quantum materials. This section reflects on the broader physical implications of the proposed model, contrasts it with conventional theories, explores potential extensions, and acknowledges inherent limitations.

7.1 Physical Implications of the Variational Inequality Approach

The NVI formulation naturally integrates physical constraints and nonlocal interactions into the modeling framework. Traditional energy minimization methods often assume smoothness or unconstrained configurations, which may not reflect realistic spin dynamics in constrained systems (Spin saturation or bounded magnetization). In contrast, the NVI approach allows for:

- Explicit incorporation of physical constraints, such as spin values constrained to specific ranges or bounded magnetic fields.
- Non-smooth transitions and metastable states, relevant in systems exhibiting hysteresis or domain walls.
- Robust handling of nonlocality, essential for materials where long-range dipolar or RKKY interactions dominate.

7.2 Comparison with Traditional Models

Several classical frameworks have been used historically to analyze phase transitions, including:

- **Landau-Ginzburg Theory:** Provides a phenomenological model using polynomial free energy expansions. However, it assumes local interactions and fails to capture spatial correlations accurately.
- **Renormalization Group (RG) Methods:** Offer deep insight into universality and critical exponents but typically rely on linear approximations or coarse-graining, making them analytically intractable in strongly correlated or spatially inhomogeneous regimes.
- **Mean-Field Approaches:** Treat spin interactions via spatial averages, neglecting fluctuations. These fail to predict critical phenomena accurately in lower dimensions or in systems with strong long-range order.

7.3 Potential Extensions

The flexibility of the NVI formulation opens avenues for future research in more exotic quantum systems:

- **Quantum Entanglement Modeling:** By enriching the state space to include entanglement measures (e.g., von Neumann entropy or concurrence), the framework could be extended to characterize *entangled spin networks*, especially in many-body localization regimes.
- **Topological Phases and Edge States:** Incorporating boundary effects and topological invariants (Chern numbers) into the energy functional could allow the study of *quantum spin Hall* or *topological insulator* phases via variational constraints.
- **Disordered Systems and Random Media:** Stochastic extensions of the variational inequality, including random coefficients or quenched disorder, could capture phenomena like *spin glass behavior* and *Anderson localization*.

7.4 Limitations and Assumptions

Despite its strengths, the proposed approach comes with several limitations:

- **Static Framework:** The current model describes equilibrium configurations. Dynamic phenomena such as relaxation, thermalization, or non-equilibrium phase transitions require time-dependent or evolution variational inequalities.
- **Continuum Approximation:** While suitable for macroscopic modeling, the continuous limit may lose fidelity for small or finite lattice systems where quantum discreteness plays a dominant role.
- **Smoothness Assumptions:** Mathematical tractability depends on the smoothness of kernels $J(x,y)$ and potentials $V(u)$. Real-world systems may exhibit singular interactions or discontinuous potentials.
- **Computational Complexity:** Long-range nonlocal interactions significantly increase the computational burden of simulation, especially in higher dimensions or fine meshes.

8. CONCLUSION

This study has introduced a novel mathematical framework based on nonlinear variational inequalities (NVIs) to model phase transitions in quantum spin systems characterized by long-range interactions. By reformulating the

governing energy equations into an NVI structure, we have successfully integrated the physical constraints, nonlocality, and nonlinear behaviors inherent in such systems. This approach not only offers a rigorous mathematical treatment of the model but also enhances its physical realism, especially in capturing sharp transitions, metastable configurations, and critical behavior.

From a mathematical perspective, the NVI framework provides a robust foundation for studying complex quantum phenomena. It leverages the power of monotone operator theory, functional analysis, and convex optimization, and allows for the derivation of existence and uniqueness results under physically relevant conditions. Additionally, the integration of long-range interaction kernels (power-law decay) within the variational formulation marks a substantial step forward in modeling spatially extended quantum systems.

On the physical side, the simulations and analysis capture key signatures of quantum phase transitions, such as bifurcations, critical thresholds, and nontrivial spin textures. Compared to traditional methods like Landau theory or mean-field approximations, this approach more accurately represents spatial correlations and the role of system constraints, offering richer insights into the ground state behavior of quantum materials with extended interactions.

Looking ahead, this work opens multiple avenues for future research:

- Higher-dimensional extensions will enable the modeling of more realistic quantum lattices, such as 2D and 3D materials, which are of great experimental interest in condensed matter physics.
- Quantum entanglement and topological phases can be incorporated into the variational structure by extending the state space and introducing additional energy terms or topological constraints.
- Non-equilibrium dynamics, including time-dependent spin evolution or driven quantum systems, could be modeled through evolutionary variational inequalities or dynamic extensions of the current framework.

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